

# Relative concentrations of carbon related defects in silicon

S.-R. G. Christopoulos<sup>1</sup> · D. C. Parfitt<sup>1</sup> · E. N. Sgourou<sup>2</sup> · C. A. Londos<sup>2</sup> · R. V. Vovk<sup>3</sup> · A. Chroneos<sup>1,4</sup>

Received: 12 May 2016/Accepted: 23 June 2016/Published online: 28 June 2016 © Springer Science+Business Media New York 2016

Abstract Irradiation induced defects in silicon are technologically important as they impact the electronic properties. Calculations based on density functional theory employing hybrid functionals have been previously used to investigate the structures and relative energies of defect clusters formed between vacancies, self-interstitials, carbon and oxygen atoms in silicon. In this study we employ a model to calculate the relative concentrations of carbon related defects in silicon. It is calculated that the carbon content has a significant impact upon the concentration of carbon-related defects. The  $C_iC_s$  defect is the most populous for all the conditions considered followed by the  $C_i$ .  $O_iSi_I$  and the  $C_iO_i$  defects.  $C_iO_iSi_I$  and the  $C_iO_i$  become increasingly important for silicon with high carbon concentrations.

## **1** Introduction

For several decades silicon (Si) has been the key material used in a range of nanoelectronic, photovoltaic, and sensor devices. It is also a model system for the understanding of point defect processes in semiconductors. The detailed understanding of

A. Chroneos ab8104@coventry.ac.uk

- <sup>1</sup> Faculty of Engineering, Environment and Computing, Coventry University, Priory Street, Coventry CV1 5FB, UK
- <sup>2</sup> Solid State Physics Section, University of Athens, Panepistimiopolis Zografos, 157 84 Athens, Greece
- <sup>3</sup> Physics Department, V. Karazin Kharkiv National University, Svobody Sq.4, Kharkiv 61077, Ukraine
- <sup>4</sup> Department of Materials, Imperial College London, London SW7 2AZ, UK

many defect processes that can impact its material properties and consequently device operation is, however, not well established [1–4]. It should be stressed that defect processes are becoming increasingly important given that the characteristic dimensions of devices are a few nanometers.

Carbon (C) is incorporated in the Si lattice during crystal growth originating from the poly-crystalline starting material, the graphitic components in the equipment and from gaseous contaminants [5–7]. C is an isovalent impurity in the Si lattice and can occupy substitutional sites  $(C_s)$ that are electrically neutral. C in the lattice leads to local strains as it is smaller than Si. C atoms interact with O<sub>i</sub> forming C–O complexes, and in particular C<sub>s</sub>–O<sub>i</sub> pairs [8]. In irradiated Si it is determined that most of the silicon interstitials  $(Si_I)$  are trapped by  $C_s$ , which are converted to C interstitials (C<sub>i</sub>) via the Watkins displacement reaction C<sub>s</sub> + Si<sub>I</sub>  $\rightarrow$  C<sub>i</sub> [9, 10]. This is important from a technological viewpoint as C<sub>i</sub> introduces in-gap states [11]. C<sub>i</sub> at room temperature can interact with O<sub>i</sub> and C<sub>s</sub> to form C<sub>i</sub>O<sub>i</sub> and  $C_iC_s$  pairs respectively [8, 12–14]. Notably, both  $C_iO_i$  and  $C_iC_s$  defects introduce [8, 12–14] states within the Si band gap, impacting devices [15-17]. This has motivated the investigation of their structure and properties [18-22]. Importantly, the  $C_i$ ,  $C_iO_i$ , and  $C_iC_s$  defects can act in the course of irradiation as nucleation centers for larger complexes such as for instance the C<sub>i</sub>O<sub>i</sub>Si<sub>I</sub> [23–27].

For completeness we need to consider the oxygen-vacancy defects. It is well established that the vacancy-oxygen pair (VO or A-center) is the most important defect in Czochralski grown Si (Cz–Si) [28]. A-centers form due to the introduction of a relatively high concentration of oxygen during the process of Czochralski crystal growth [28]. The oxygen interstitials (O<sub>i</sub>) associate with vacancies (V) to form VO pairs via the reaction  $V + O_i \rightarrow VO$  [28–30]. These VO pairs are technologically important as they are both electrically and optically active and therefore it may be necessary to control their formation and suppress their concentration [31, 32].

Defect clusters compete for point defects, and in that respect, we need to consider the formation of all defects and calculate their relative populations [33]. The aim of the present study is to calculate the relative concentrations of carbonrelated defects within the mass action analysis framework.

## 2 Methodology

#### 2.1 Definitions

Experiment can provide information on the thermal stability and formation of defect clusters and in that respect DFT can act synergistically through the calculation of the binding energies,  $E_b$ , defined by:

$$E_b = E_{\text{defectcluster}} - \Sigma E_{\text{isolateddefects}} \tag{1}$$

With this definition negative binding energies imply that a cluster is energetically favourable with respect to its constituent isolated components.

The concentration of an XY complex [XY] with a binding energy  $E_b$  relative to the concentration of unbound defects X and Y (i.e. [X] and [Y] respectively) is given by:

$$\frac{[XY]}{[X][Y]} = \exp\left(-E_{b/k_{B}T}\right) \tag{2}$$

where  $k_B$  is Boltzmann's constant and T is the absolute temperature.

Mass action techniques were previously used to calculate the relative concentrations of donor-vacancy clusters and the impact of codoping in germanium (Ge) [34, 35] and hydrogen related defects in amorphous Si [36]. The approach is predictive as it resulted in the first report of donor atom-vacancy clusters in Ge and the formation of fluorine clusters in Ge both of which were thereafter confirmed by experimental studies [35, 37]. In this approach the temperature dependence of defect concentrations at equilibrium conditions can be calculated. Nevertheless, kinetics may play a role, as will non-equilibrium conditions (i.e. irradiation). In the present work, the relative concentrations of carbon related defects are calculated under equilibrium conditions.

#### **3** Results and discussion

## 3.1 Structure of defects

From an experimental viewpoint, the main defect species identified by FTIR spectroscopy are C<sub>i</sub>, C<sub>i</sub>O<sub>i</sub>, C<sub>i</sub>C<sub>s</sub>, C<sub>i</sub>O<sub>i</sub>Si<sub>I</sub>,

VO,  $V_2$ , and  $VO_2$ , (Ref. 38 and references therein). The relations for their formation are:

$$C_s + Si_I \rightarrow C_i$$
 (3)

$$C_i + O_i \rightarrow C_i O_i$$
 (4)

$$C_i + C_s \to C_i C_s \tag{5}$$

$$C_i O_i + Si_I \to C_i O_i Si_I \tag{6}$$

$$V + O \rightarrow VO$$
 (7)

$$V + V \to VV$$
 (8)

$$VO + O_i \rightarrow VO_2$$
 (9)

The structure of  $C_i C_s$  defects (known as G-centers) have been extensively investigated using DFT calculations [39, 40]. There are three structures of defects (refer to Fig. 1 of Ref. 40). To study the structure of the successive members of the CiOi(SiI)n family of defects it is necessary to begin with the C<sub>i</sub>O<sub>i</sub> defect. The C<sub>i</sub>O<sub>i</sub> defect has two forms [19] the O-form (refer to Fig. 1a) and the R-form (refer to Fig. 1b) [19]. R denotes the ring formed by the O<sub>i</sub>,  $C_i$  and two Si lattice atoms. The  $C_iO_i(Si_I)$  defect forms by the attachment of a Si<sub>I</sub> atom in the C<sub>i</sub>O<sub>i</sub> defect (refer to Eq. 6). There are at least three configurations of the C<sub>i</sub>O<sub>i</sub>(Si<sub>I</sub>) defect that have been previously investigated (refer to Fig. 1c-e) [20, 25, 41]. The first two configurations of the  $C_iO_i(Si_1)$  can be formed by the two configurations of the C<sub>i</sub>O<sub>i</sub> pair (Fig. 1a,b) by the addition if a Si<sub>I</sub>. The third configuration of the  $C_iO_i(Si_I)$  (Fig. 1e), is the structure proposed by Estreicher and Backlund [25, 42] using DFT. In this structure the O and C atoms are again bonded in an R-form. Upon annealing an additional Si<sub>I</sub> atom can be captured by the CiOi(SiI) defect to form a  $C_iO_i(Si_I)_2$  structure, however, these larger defects are beyond the scope of the present study. Further DFT studies will be required to clarify the structure and energetics of these larger defects. The structure of VO,  $V_2$ , and  $VO_2$ defects is reported in previous studies [43].

#### 3.2 Model

To calculate the relative concentrations of the defects described by Eqs. 3–9, we have to solve simultaneously the following relations:

$$\frac{[C_i]}{[C_s][Si_I]} = \exp\left(\frac{2.34eV}{k_BT}\right) \tag{10}$$

$$\frac{[C_i O_i]}{[C_i][O_i]} = \exp\left(\frac{1.60eV}{k_BT}\right)$$
(11)

**Fig. 1** Schematic representations of the C<sub>i</sub>O<sub>i</sub> (**a**, **b**) and C<sub>i</sub>O<sub>i</sub>Si<sub>I</sub> (**c**–**e**) defects



$$\frac{[C_i C_s]}{[C_i][C_s]} = \exp\left(1.68eV/k_BT\right)$$
(12)

$$\frac{[C_i O_i S i_I]}{[C_i][O_i][S i_I]} = \exp\left(3.14eV_{k_BT}\right)$$
(13)

$$\frac{[VV]}{[V][V]} = \exp\left(2.52eV_{k_BT}\right) \tag{14}$$

$$\frac{[VO_i]}{[V][O_i]} = \exp\left(\frac{2.21eV}{k_BT}\right)$$
(15)

$$\frac{[VO_2]}{[V][O_i][O_i]} = \exp\left(4.05eV_{k_BT}\right)$$
(16)

This is the mass action framework and can provide useful information within the approximations described above. Additionally, it is necessary to have accurate binding energies. The present binding energies in these relations where either derived by Wang et al. [26, 27]. or by using the same state of the art hybrid functional DFT method. The thermal equilibrium concentration of Si<sub>I</sub> and *V* is very low due to their relatively high formation energies. It is assumed here that under irradiation we must account for the formation of a highly athermal concentration of these defects which may interact and stabilize oxygen and carbon complexes. We consider here the impact of initial carbon related defects.

#### 3.3 Impact of carbon concentration

In Fig. 2 the dependence of carbon related defects with respect to the initial carbon content for two characteristic temperatures (T = 500 and 1000 K) is given. For both

temperatures considered it can be seen that the carbonrelated defects are dominated by  $C_i$  and  $C_iC_s$ . The latter defect becomes more populous when the initial carbon concentration exceeds  $3 \times 10^{18}$  cm<sup>-3</sup> (refer to Fig. 2).

As it is expected the increase of the C concentration has a significant impact on the concentration of carbon-related defects. The concentrations of defect clusters considered is significantly increased with the increase of temperature. This is more profound for  $C_iO_i$  and  $C_iO_iSi_I$ . Notably,  $C_iC_s$  can be beneficial as it can be used to improve the Si optical emitters [17, 38], whereas  $C_iO_i$  and  $C_iO_iSi_I$  can have a deleterious impact. Therefore we propose here that processing should be limited to lower temperatures to allow the formation of  $C_iC_s$  but avoid the formation of  $C_iO_i$  and  $C_iO_iSi_I$  defects.

In the present study it is deduced that the carbon concentration has a limited impact on the concentration of the main oxygen-vacancy defects in Si. This is consistent to a recent model that considered the effect of isovalent doping (Ge, Sn, Pb) in Si [44]. Although carbon was not considered in that study it was demonstrated that only defects that had a strong interaction with vacancies (i.e. Sn and Pb) had a significant impact on the VO and VO<sub>2</sub> concentrations (refer to Fig. 4 of Ref. 44). Carbon has only a small interaction with vacancies and additionally the  $C_iO_i$  defect has a lower binding energy as compared to VO, therefore it is expected that carbon will have a limited impact upon the concentration of the oxygen-vacancy defects [38, 42].

## 3.4 Other considerations

It should be stressed that the present study involves the calculation of relative defect concentrations at equilibrium

Fig. 2 The dependence of carbon related defects with respect to the initial unbound carbon concentration for T = 500 K and T = 1000 Kand initial  $[O_i] = [Si_I] = [V] = 10^{18}$ cm

be incorporated in future work.



conditions aiming to illustrate the impact of binding energies (i.e. mainly local strain relaxation). Most experiments typically involve irradiation and are under nonequilibrium conditions. Under these conditions kinetics will play a role as will extended defects and external sinks (i.e. surfaces). These effects are being considered and will

Si<sub>1-x-v</sub>Ge<sub>x</sub>Sn<sub>v</sub> binary and ternary alloys are promising materials for numerous applications including detectors, photovoltaics, lasers and microelectronic devices [45, 46]. Although, mass action models have been used to other semiconducting materials such as Ge and Si1-xGex it becomes increasingly difficult to model ternary materials [47–49] Considering the kinetics of the processes the picture becomes even more complicated in ternary semiconductors but forms a fruitful field for further experimental and theoretical investigations [50-53].

#### 4 Summary

We present a model to describe carbon related defects in Si. The input parameters of the model are state of the art hybrid DFT binding energies. The model concerns equilibrium conditions and calculates the relative concentration of carbon related defects for a range of carbon concentrations.  $C_i$  and  $C_iC_s$  are the dominant carbon related defects, whereas CiOiSiI and CiOi defects become increasingly important for high temperatures and the higher initial carbon concentrations. It was previously calculated that the oversized (Pb and Sn) isovalent dopants impact the oxygen-vacancy defects [44], whereas here we showed that carbon does not. This is due to the limited formation of  $C_iO_i$  defects and the insignificant interaction of carbon atoms with vacancies. Therefore, carbon does not associate with both componentes of the oxygen-vacancy defects.

Initial Conc. ([C] $_{init.}$ ), cm<sup>-3</sup>

Acknowledgments SRGC and AC are grateful for funding from the Lloyd's Register Foundation, a charitable foundation helping to protect life and property by supporting engineering-related education, public engagement and the application of research.

#### References

- 1. V.V. Emtsev, N.V. Abrosimov, V.V. Kozlovskii, G. Oganesyan, Semiconductors 48, 1438 (2014)
- 2. C.A. Londos, E.N. Sgourou, D. Timerkaeva, A. Chroneos, P. Pochet, V.V. Emtsev, J. Appl. Phys. 114, 113504 (2013)
- 3. V.V. Voronkov, R. Falster, C.A. Londos, E.N. Sgourou, A. Andrianakis, H. Ohyama, J. Appl. Phys. 110, 093510 (2011)
- 4. A. Chroneos, C.A. Londos, E.N. Sgourou, J. Appl. Phys. 110, 093507 (2011)
- 5. B.O. Kolbersen, A. Muhlbauer, Solid State Electron. 25, 759 (1982)
- 6. R.C. Newman, Mat. Res. Soc. Symp. Proc. 59, 403 (1986)
- 7. W. Scorupa, R.A. Yankov, Mater. Chem. Phys. 44, 101 (1996)
- 8. G. Davies, E.C. Lightowlers, R.C. Newman, A.S. Oates, Semicond. Sci. Technol. 2, 524 (1987)
- 9. R.C. Newman, A.R. Bean, Radiat. Eff. 8, 189 (1970)
- 10. G.D. Watkins, K.L. Brower, Phys. Rev. Lett. 36, 1329 (1976)
- 11. C.A. Londos, Semicond. Sci. Technol. 5, 645 (1990)
- 12. E.V. Lavrov, L. Hoffmann, B.B. Nielsen, Phys. Rev. B 60, 8081 (1999)
- 13. C.A. Londos, M.S. Potsidi, E. Stakakis, Phys. B 340-342, 551 (2003)
- 14. C.-L. Liu, W. Windl, L. Borucki, S. Lu, X.-Y. Liu, Appl. Phys. Lett. 80, 52 (2002)
- 15. G. Davies, R.C. Newman, in Handbook of Semiconductors, vol. 3, ed. by S. Mahajan (Elsevier, Amsterdam, 1994), p. 1557
- 16. M. Yamaguchi, A. Khan, S.J. Taylor, K. Ando, T. Yamaguchi, S. Matsuda, T. Aburaya, J. Appl. Phys. 86, 217 (1999)
- 17. K. Murata, Y. Yasutake, K. Nittoh, S. Fukatsu, K. Miki, AIP Adv. 1, 032125 (2011)
- 18. J.M. Trombetta, G.D. Watkins, Appl. Phys. Lett. 51, 1103 (1987)
- 19. R. Jones, S. Oberg, Phys. Rev. Lett. 68, 86 (1991)
- 20. J. Coutinho, R. Jones, P.R. Briddon, S. Oberg, L.I. Murin, V.P. Markevich, J.L. Lindstrom, Phys. Rev. B 65, 014109 (2001)
- 21. P. Leery, R. Jones, S. Oberg, V.J.B. Torres, Phys. Rev. B 55, 2188 (1997)

- R.B. Capaz, A. Dal Pino Jr., J.D. Joannopoulos, Phys. Rev. B 58, 9845 (1998)
- A. Mattoni, F. Bernantini, L. Colombo, Phys. Rev. B 66, 195214 (2002)
- 24. C.A. Londos, M.S. Potsidi, G.D. Antonaras, A. Andrianakis, Phys. B 376–377, 165 (2006)
- 25. D.J. Backlund, S.K. Estreicher, Phys. Rev. B 77, 205205 (2008)
- H. Wang, A. Chroneos, C.A. Londos, E.N. Sgourou, U. Schwingenschlögl, Appl. Phys. Lett. 103, 052101 (2013)
- H. Wang, A. Chroneos, C.A. Londos, E.N. Sgourou, U. Schwingenschlögl, Sci. Rep. 4, 4909 (2014)
- R.C. Newman, R. Jones, Oxygen in silicon, in *Semiconductors and Semimetals*, vol. 42, ed. by F. Shimura (Academic Press, Orlando, 1994), p. 289
- J.W. Corbett, G.D. Watkins, R.S. McDonald, Phys. Rev. A 135, 1381 (1964)
- 30. C.A. Londos, Phys. Status Solidi. A 102, 639 (1987)
- C.A. Londos, E.N. Sgourou, D. Hall, A. Chroneos, J. Mater. Sci.: Mater. Electron. 25, 2395 (2014)
- D.J. Hall, N.J. Murray, J.P.D. Gow, D. Wood, A. Holland, J. Instrum. 9, C12004 (2014)
- F.A. Kröger, V.J. Vink, in *Solid State Physics*, vol. 3, ed. by F. Seitz. D. Turnbull (Academic, New York, 1956), p. 307
- A. Chroneos, R.W. Grimes, B.P. Uberuaga, S. Brotzmann, H. Bracht, Appl. Phys. Lett. 91, 192106 (2007)
- G. Impellizzeri, S. Boninelli, F. Priolo, E. Napolitani, C. Spinella, A. Chroneos, H. Bracht, J. Appl. Phys. 109, 113527 (2011)
- 36. M.J. Powell, S.C. Deane, Phys. Rev. B 53, 10121 (1995)
- D.J. Sprouster, C. Campbell, S.J. Buckman, G. Impellizzeri, E. Napolitani, S. Ruffell, J.P. Sullivan, J. Phys. D Appl. Phys. 46, 505210 (2013)
- A. Chroneos, E.N. Sgourou, C.A. Londos, U. Schwingenschlögl, Appl. Phys. Rev. 2, 021306 (2015)

- F. Zirkelbach, B. Stritzker, K. Nordlund, J.K.N. Linder, W.G. Schmidt, E. Rauls, Phys. Rev. B 84, 064126 (2011)
- H. Wang, A. Chroneos, C.A. Londos, E.N. Sgourou, U. Schwingenschlögl, J. Appl. Phys. 115, 183509 (2014)
- G. Davies, S. Hayama, S. Hao, B.B. Nielsen, J. Coutinho, M. Sanati, S.K. Estreicher, K.M. Itoh, Phys. Rev. B 71, 115212 (2005)
- 42. D.J. Backlund, S.K. Estreicher, Phys. B 401-402, 163 (2007)
- C.A. Londos, L.G. Fytros, G.J. Georgiou, Defect Diffus. Forum 171–172, 1 (1999)
- S.R.G. Christopoulos, D.C. Parfitt, E.N. Sgourou, C.A. Londos, A. Chroneos, J. Mater. Sci.: Mater. Electron. 27, 4385 (2016)
- A.V.G. Chizmeshya, C. Ritter, J. Tolle, C. Cook, J. Menendez, J. Kouvetakis, Chem. Mater. 18, 6266 (2006)
- A. Chroneos, C. Jiang, R.W. Grimes, U. Schwingenschlögl, H. Bracht, Appl. Phys. Lett. 95, 112101 (2009)
- A. Chroneos, H. Bracht, R.W. Grimes, B.P. Uberuaga, Mater. Sci. Eng., B 154–155, 72 (2008)
- 48. A. Chroneos, H. Bracht, Appl. Phys. Rev. 1, 011301 (2014)
- 49. A. Chroneos, R.W. Grimes, H. Bracht, J. Appl. Phys. 106, 063707 (2009)
- E.N. Sgourou, D. Timerkaeva, C.A. Londos, D. Aliprantis, A. Chroneos, D. Caliste, P. Pochet, J. Appl. Phys. 113, 113506 (2013)
- G. Brenet, D. Timerkaeva, E.N. Sgourou, C.A. Londos, D. Caliste, P. Pochet, J. Appl. Phys. 118, 125706 (2015)
- A. Chroneos, C. Jiang, R.W. Grimes, U. Schwingenschlogl, H. Bracht, Appl. Phys. Lett. 94, 252104 (2009)
- S.T. Murphy, A. Chroneos, C. Jiang, U. Schwingenschlogl, R.W. Grimes, Phys. Rev. B 82, 073201 (2010)